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1 **Identifying Cold Pool scales over Complex Topography**  
2 **using TaiwanVVM simulations**

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## Abstract

24

25 In this study, our objective is to identify the appropriate cold pool scales over Taiwan's  
26 complex topography during predominant afternoon thunderstorms under a local-circulation  
27 dominated weather regime in summer. We utilize semi-realistic TaiwanVVM simulations,  
28 which cover the entire area of Taiwan, to investigate this phenomenon. Our findings reveal  
29 that when buoyancy is defined using a conventional environmental scale (109 km), the cold  
30 pool locations do not align with the precipitation areas, instead being concentrated mainly  
31 along the mountain ridges. We hypothesize that this discrepancy arises from the  
32 environmental scale at which cold pool buoyancy operates. To assess this, we conducted  
33 systematic analyses and the results show that an optimal environmental scale of  
34 approximately 7 to 11 km (about 3 times of the 75<sup>th</sup> and 90<sup>th</sup> percentile of the precipitation  
35 object length) can be identified. The statistics of cold pool frequency better align with  
36 precipitation hotspots, characterized by evaporative cooling over the plains and increased  
37 water loading within the core of precipitation objects over the mountains. We demonstrate  
38 that this method effectively captures the shift in cold pools associated with precipitation  
39 responses in a warming climate in Taiwan. This work highlights the importance of using an  
40 appropriate environmental scale when estimating buoyancy over complex topography.

41 **Keywords** cold pool; complex topography; buoyancy

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43 **1. Introduction**

44 Diurnal convection over tropical islands is well known for its significant contribution to  
45 overall precipitation (Nesbitt and Zipser, 2003; Yang and Slingo, 2001). This convection is  
46 typically thermally driven and produces intense, short-duration heavy rainfall (Romatschke  
47 and Houze, 2011; X. Chen et al., 2015; Song and Zhang, 2020; Jian et al., 2021; P.-J. Chen  
48 et al., 2021; Krishna et al., 2021). In Taiwan, particularly over its complex topography, diurnal  
49 convection primarily occurs in the mountainous regions of southwest Taiwan, accounting for  
50 up to 40% of summer diurnal rainfall under local circulation-dominant conditions (Chang et  
51 al., 2023).

52 During convection development, the cold pool, driven by cold and dry downbursts from  
53 the cooling effect of precipitation evaporation (Doswell, 2001; Benjamin, 1968; Von Karman,  
54 1940; Keulegan, 1957; Kuenen, 1950; Byers and Braham, 1949; Webster and Lukas, 1992),  
55 induces mechanical lifting for new convection initiation from the leading gust front (cold pool  
56 edge) (Droegemeier and Wilhelmson, 1985; Thorpe et al., 1982). The cold pool also triggers  
57 convection initiation from the surrounding cold pool under unstable atmospheric conditions  
58 (Tompkins, 2001a, 2001b). Moreover, convection initiation is more likely induced by cold  
59 pool collisions than by isolated cold pool lifting (Feng et al., 2015). Additionally, the collision  
60 of cold pools generated by convection with the sea-breeze front triggers new convection in  
61 the Taipei Basin (Kuo and Wu, 2019; Miao and Yang, 2020, 2022). This process  
62 demonstrates how cold pool dynamics extend convection development and influence its  
63 spatial distribution on tropical islands.

64 There are three primary methods used to identify cold pools. The buoyancy method  
 65 involves the vertical integration of negative buoyancy derived from the density potential  
 66 temperature ( $\theta_v$ ) anomaly (Schlemmer and Hohenegger, 2015; Tompkins, 2001a; Feng et  
 67 al., 2015). The radial derivatives method identifies the cold pool edge by locating the  
 68 steepest gradient of buoyancy (Drager and van den Heever, 2017). The radial gradient-  
 69 based method identifies the cold pool edge where strong and narrow convergence regions  
 70 exist near the surface (Fournier and Haerter, 2019; Henneberg et al., 2020).

71 The radial derivatives and radial gradient-based methods primarily identify the cold pool  
 72 edge without considering the cold pool structure. In the radial gradient-based method, the  
 73 surface convergence region calculated from near-surface wind is notably affected by steep  
 74 terrain. Conversely, the buoyancy method not only identifies the cold pool edge but also  
 75 characterizes the cold pool structure. This method can be used to explore the contribution  
 76 of cold pools to the buoyancy structure near convective regions, which is related to the non-  
 77 local effect in convection (Kuo and Neelin, 2022), particularly in diurnal convection over  
 78 complex topography in Taiwan.

79 In our study, we quantify cold pools based on cold pool intensity ( $\text{m s}^{-1}$ ) using the  
 80 buoyancy method (Tompkins, 2001a; Feng et al., 2015). The buoyancy B is given by:

$$81 \quad B = \frac{g(\theta_v - \bar{\theta}_v)}{\theta_{v_0}} \quad (\text{Eqn. 1})$$

$$82 \quad \theta_v = \theta(1 + 0.608q_v - q_l) \quad (\text{Eqn. 2})$$

$$83 \quad \text{cold pool intensity} = \sqrt{-2 \int_0^h B dz} \quad (\text{Eqn. 3})$$

84 where  $B$  is the buoyancy ( $\text{m s}^{-2}$ ),  $g$  is the acceleration due to gravity ( $\text{m s}^{-2}$ ),  $\theta$  is the potential  
 85 temperature,  $q_v$  is the mass mixing ratio of water vapor, and  $q_l$  includes cloud condensate  
 86 water, rainwater, and ice. The  $\theta_{v0}$  represents the reference state (Stevens, 2005, Yano et  
 87 al., 2004). The  $\bar{\theta}_v$  represents the  $\theta_v$  average under conventional environmental scale  
 88 (109km, the size of red box region in Fig.4a), which is used in previous studies (Feng et al.,  
 89 2015)

90 The cold pool intensity is determined by integrating  $B$  from the point where it first exceeds  
 91  $-0.003$  ( $\text{m s}^{-2}$ ) (following Feng et al., 2015), which helps to estimate the cold pool depth ( $h$ )  
 92 above the land surface. Fig. 1 shows the cold pool intensity in plain and mountainous regions  
 93 from TaiwanVVM. However, the cold pool intensity is notably stronger and primarily  
 94 concentrated near mountain ridges around 500 meters high, while it is nearly  $0 \text{ m s}^{-1}$  in the  
 95 plains. The appropriate cold pool intensity for a tropical island is approximately  $4 \text{ m s}^{-1}$   
 96 (Drager & van den Heever, 2017; Zuidema et al., 2017).

97 In this study, we use the precipitation object length to determine the environmental scale,  
 98 based on the concept that cold pools primarily occur near precipitation regions generated  
 99 by convection. As the precipitation area increases, the area with cold pools also increases,  
 100 indicating that the environmental field that can contain the negative buoyancy of the cold  
 101 pool generated by convection also expands. Therefore, by using the precipitation area and  
 102 expanding outward by a certain proportion as the environmental field for the negative

103 buoyancy of the cold pool, we can identify the signals of cold pools generated near the  
104 convection. When the environmental scale is small, the resulting cold pools may be weak  
105 and fragmented, whereas when the environmental scale is large, the buoyancy calculation  
106 could be dominated by terrain surface heating/cooling or mountain-valley wind effects. In  
107 this study, we will identify an appropriate environmental scale for cold pools driven by  
108 convective precipitation over complex topography in Taiwan.

109 To investigate the relationship between cold pools and convection in complex terrain, we  
110 will calculate the statistics of diurnal precipitation objects. Cold pools formed by  
111 precipitation can trigger new convection, leading to convection merging process and  
112 stronger rainfall. Thus, cold pools not only initiate new convection but also act as a bridge in  
113 the process of convection merging, which is the primary scenario we studied. For large  
114 convection scales, the situation becomes more complex; for example, the lower atmosphere,  
115 influenced by heavy precipitation, becomes nearly saturated, and cold pools primarily form  
116 within the core of convective precipitation. Such scenario is beyond the scope of this study.

117 This study is organized as follows. Section 2 provides the model description and initial  
118 conditions of the TaiwanVVM ensemble member, and details of the Identification of the  
119 precipitation object length. Section 3 first presents the patterns of precipitation evolution in  
120 Taiwan under a local-circulation dominated weather regime, then examines the behavior of  
121 cold pool calculated from environmental scale structure and at last discuss the contribution  
122 on cold pool over complex topography in Taiwan. Section 4 discusses cold pool hotspot

123 changes under Pseudo-Global Warming (PGW) scenario (Chen et al., 2024). The summary  
124 and conclusions are discussed in Section 5.

## 125 **2. Data and Method**

### 126 **2.1 Model and experiment description**

127 In this study, we use 30 cases of TaiwanVVM semi-realistic simulations as the member  
128 ensemble representing both the current and PGW climate, following Chang et al. (2021) and  
129 Chen et al. (2024). In the semi-realistic, the *realistic* aspect refers to the use of real  
130 topography and land-type distributions to construct the boundary conditions in the model to  
131 represent Taiwan island. Meanwhile, the *semi* aspect indicates that the atmospheric initial  
132 conditions are based on the 00Z Banqiao sounding data, assuming a horizontally  
133 homogeneous initial state which is typically used in Large-Eddy Simulations (LES). This  
134 setup focuses on the development of local circulation and convection-driven precipitation  
135 processes through surface boundary perturbations and differential heating caused by the  
136 terrain and land surface. TaiwanVVM (Wu et al., 2019) predicts the three-dimensional  
137 anelastic vorticity while the vertical velocity is diagnosed by solving a 3-D elliptic equation.  
138 The model uses predicted particle properties microphysics scheme (P3, Morrison and  
139 Milbrandt, 2015) and use the Rapid Radiative Transfer Model for GCMs (RRTMG; Iacono  
140 et al., 2008) as the radiative parameterization. To estimate the surface fluxes, the model  
141 uses the flux–profile relationship of Deardorff (1972), and the eddy viscosity and diffusivity  
142 coefficients depending on deformation and stability (Shutts and Gray, 1994) as the first-order



143 turbulence closure. TaiwanVVM includes realistic Taiwan orography (Wu and Arakawa,  
144 2011) at 500m horizontal resolution and coupling Noah land surface model (Chen and  
145 Dudhia, 2001; Chen et al., 1996) version 3.4.1 to capture the differential surface heating  
146 over complex topography as the boundary conditions.

147 We adopted the definition of local-circulation dominated weather regime for the summer  
148 months (May to September) from 2005 to 2015 and selected 30 sounding dates as  
149 representative cases from Chang et al. (2021) as our TaiwanVVM simulation. The initial  
150 thermodynamic soundings (Fig. 2) represent a conditionally unstable environment typical of  
151 Taiwan's summer, which is the only difference among the simulations. This setup is to  
152 highlight the interactions between physical processes that control the development of local  
153 circulation and convection, as well as to examine the sensitivity of these processes to key  
154 environmental factors (Chen et al., 2024). The model setting of the whole cases we studied  
155 is in Table 1. For the discussion on cold pool frequency change, the PGW scenario is also  
156 considered.

## 157 **2.2 Identification of the precipitation object length and environmental scale**

158 The precipitation object length is determined by the precipitation object, which is  
159 connecting the raining grid points ( $\geq 5 \text{ mm h}^{-1}$ ) (Fig. 3a) in a two-dimensional (x-y) manner,  
160 following Tsai and Wu (2017). The precipitation object length (Fig. 3b) is calculated from the  
161 square root of the precipitation object's size, which also represents the convective scale in  
162 mountain region in Taiwan.

$$163 \quad B = \frac{g(\theta_v - \tilde{\theta}_v)}{\theta_{v0}} \quad (\text{Eqn.4})$$

164 Next, eqn. (1) is modified from eqn. (4), where we  $\bar{\theta}_v$  is  $\tilde{\theta}_v$ . Note that  $\bar{\theta}_v$  is evaluated  
 165 for the environmental scale of 109 km (red box region in Fig.4a), but  $\tilde{\theta}_v$  is evaluated by  
 166 using environmental scales that will be derived later. The  $\tilde{\theta}_v$  is the horizontal average of  $\theta_v$   
 167 calculated for the square domain with the edge length of an environmental scale. For  
 168 example, the black box in Fig.3b which covers the precipitation object, is selected to  
 169 calculate the local environmental state for convection developing under complex topography.  
 170 In this study, we will evaluate the environment scale based on the precipitation object length.  
 171 In each ensemble member, the precipitation object length is calculated based on four  
 172 specific percentiles (25<sup>th</sup>, 50<sup>th</sup>, 75<sup>th</sup>, and 90<sup>th</sup>) (We will illustrate the meaning of choosing  
 173 these percentiles in section 3.1) for each time period, reflecting the evolution of precipitation  
 174 object length within the studied region. The selection of these percentiles aims to identify  
 175 the precipitation object lengths at each time period in Southwest Taiwan.

## 176 **3. Results**

### 177 **3.1 Overall Results**

178 To understand the spatial distribution of precipitation object over southern Taiwan, we  
 179 count the precipitation occurrence frequency where the rain rate is over 5 mm h<sup>-1</sup>. The result  
 180 shows that precipitation mainly occurs in the southwestern mountain hill (500-1500 m) and  
 181 high mountain (>1500 m) regions (Fig. 4a). Within this area, the highest precipitation  
 182 frequency is primarily found in the mountain hill region, with a frequency of about 10-15%,

183 indicating a precipitation hotspot. Based on this result, the following study focuses on the  
184 terrestrial region inside the red box in Fig. 4a. We hypothesize that the precipitation hotspots  
185 are caused by sea-breeze inflow in the plain region (<500 m) and mountain winds from  
186 surface heating or downbursts driven by convection in the high mountain region. To verify  
187 this hypothesis, we analyze the mean precipitation evolution in three studied terrains: plains,  
188 mountain hills, and high mountains.

189 To further understand when the precipitation mainly occurred, we analyze the  
190 precipitation evolution in plain, mountain hill and high mountain regions. The result shows  
191 that the time offset of the mean precipitation peak reflects the convection development in  
192 these area (Fig. 4b). The earliest and most intense precipitation peak among the three areas  
193 occurs in the high mountain region ( $2.1 \text{ mm h}^{-1}$ ) around 14 LST, and it decreases until 21  
194 LST. In the plain region, the mean precipitation slightly increases by  $0.1 \text{ mm h}^{-1}$  at 16 LST,  
195 likely due to the convective system developing from the mountain hill region, where the  
196 mean precipitation around 16 LST is still increasing.

197 In the plain region, precipitation before the diurnal peak at 15:30 LST, with a diurnal  
198 amplitude of about  $0.5 \text{ mm h}^{-1}$ , which is much smaller than in other regions. In the mountain  
199 hill region, the diurnal peak of mean precipitation is  $1.9 \text{ mm h}^{-1}$  around 16 LST, which is  
200 delayed by about 2 hours compared to the high mountain region and 0.5 hours compared  
201 to the plain region. We infer that the additional forcing comes from mountain winds and  
202 outflows from the high mountain region, as well as the sea-breeze inflow from the plain

203 region. After 16 LST, the mean precipitation in the mountain hill region was  $0.25 \text{ mm h}^{-1}$  and  
204  $0.5\text{-}1.5 \text{ mm h}^{-1}$  greater than in the high mountain and plain regions, respectively. The mean  
205 precipitation remains higher in the mountain hill region than in the other two regions until  
206 night, with less surface heating from 17 to 21 LST. This highlights the persistence of  
207 convection and the precipitation hotspot (Fig. 4a), which we infer is due not only to the  
208 additional forcing from the high mountain and plain regions but also to the interactions of  
209 local convection with cold pools and surface inflow over the complex topography,  
210 maintaining the precipitation events in the mountain hill region.

211 From Fig. 4c, the result shows that the diurnal peak of largest 90<sup>th</sup> precipitation object  
212 length is about 5.5 km, which reveals that the scale of the afternoon thunderstorm in Taiwan  
213 is small. To identify the location of such precipitation objects, we label the most frequent  
214 occurrence altitude among ensemble members using symbols: circles for below 500 m,  
215 squares for 500–1500 m, and triangles for above 1500 m, respectively. The square on the  
216 lines shows that the precipitation object, which the convection mainly occurs in mountain hill  
217 region after 14 LST until night (21 LST). In 75<sup>th</sup> to 90<sup>th</sup> percentile, the diurnal peak is about  
218 3.2 to 5.5 km, and the diurnal variation is larger (about 1.7 to 3.5 km) than in 25<sup>th</sup> and 50<sup>th</sup>  
219 (smaller than 0.5 km), means that the object length in 75<sup>th</sup> and 90<sup>th</sup> percentile relatively  
220 reflects the diurnal variation of convection. The analysis of precipitation object length reveals  
221 that the scale of afternoon thunderstorms in Taiwan is generally small, even at the 90<sup>th</sup>  
222 percentile. This suggests that, due to constraints imposed by the complex topography, the

223 environmental scale of cold pools must be larger than the precipitation object but small  
224 enough to avoid interference from the terrain. Therefore, we will design systematic analyses  
225 based on the length scales obtained in this section as a guide to evaluate how to select an  
226 appropriate environmental scale over complex terrain.

### 227 **3.2 The analysis of environmental scale for cold pool**

228 Based on the precipitation object analysis in Section 3.1, we will test the environmental  
229 scale using the diurnally averaged 25<sup>th</sup>, 50<sup>th</sup>, 75<sup>th</sup>, and 90<sup>th</sup> percentiles of precipitation object  
230 lengths, which are 0.88, 1.48, 2.42, and 3.67 km, respectively. By extending these spatial  
231 scales outward by an additional factor of the object length on all sides, we select three times  
232 the original scale, which is large enough to cover the precipitation objects, resulting in 2.66,  
233 4.44, 7.25, and 11.01 km, respectively. This approach will help determine the lower bound  
234 of the environmental scale. Additionally, we extend the scale by incorporating larger  
235 environmental scales—33 km, 66 km, 89 km and 100 km—to determine the upper bound,  
236 which is constrained by the spatial scale of the terrain. This methodology allows us to  
237 evaluate the environmental scale that can be applied to other tropical islands with complex  
238 topography.

239 It is important to note that cold pool calculations will be excluded if the ratio of missing  
240 grid points (e.g., due to terrain) to total grid points within the environmental scale domain  
241 exceeds 0.5 for a specific land surface point (i.e., the black point in Fig. 3b).

242 To clearly illustrate the contrast in cold pool distribution at different scales, Fig. 5 presents

243 results for four selected scales (2.66 km, 7.25 km, 33 km, and 89 km, respectively) and  
244 focuses on the plain and mountain regions in one of the ensemble members. The  
245 environmental scale is represented as a range, as the results between the 25<sup>th</sup> and 50<sup>th</sup>  
246 percentiles are similar, while those between the 75<sup>th</sup> and 90<sup>th</sup> percentiles are also  
247 comparable. For simplicity, these ranges are referred to as 2–4 km and 7–11 km,  
248 respectively. Similarly, the results between 33 km and 60 km are comparable, as are those  
249 between 89 km and 100 km. For simplicity, these ranges are referred to as 33–60 km and  
250 89–100 km, respectively.

251 Figs. 5a and b show that with a small environmental scale (2–4 km), the cold pool  
252 distribution shows weak signals surrounding the precipitation region in both plain and  
253 mountainous areas. Even though the cold pool distribution exhibits a gust-front-like structure,  
254 it is severely fragmented and appears as only a surface perturbation in the precipitation  
255 region. This indicates that the scale is too small, capturing only grid-point scale negative  
256 buoyancy signals and failing to show the cold pool structure driven by the local convective  
257 system.

258 With a larger environmental scale of approximately 7–11 km, the cold pool structure  
259 forms a well-defined ring structure with an appropriate cold pool intensity of approximately  
260  $4 \text{ m s}^{-1}$  in the plain region (Fig. 5c). In the mountain region (Fig. 5d), where the terrain is  
261 divided, the cold pool structure from large convective systems cannot form a complete ring  
262 structure. Currently, the main focus is on observing the cold pools generated by convection

263 before the convective cells merge and grow, which subsequently trigger new convection and  
264 aid in further process of convection merging. The identified environmental scale ranges in  
265 7–11 km, and the calculated cold pool signals mainly reflect the precipitation objects that  
266 contain a single stronger precipitation core. In precipitation objects larger than 20 km, even  
267 when using an environmental scale of 7–11 km to calculate the cold pool, it reflects the cold  
268 pool signals generated around the stronger precipitation cores within the precipitation  
269 objects.

270 When the environmental scale increases to 33-60 km, the cold pool intensity becomes  
271 excessively strong, and the structure starts to become vague and mismatches with the  
272 precipitation region in the plain region (Fig. 5e). In the mountain region (Fig. 5f), although  
273 some of the cold pool distribution concentrates in the precipitation region similar to Fig. 5d,  
274 it follows the mountain ridge without aligning with the precipitation region. Furthermore, the  
275 cold pools within strong precipitation areas are primarily distributed within larger precipitation  
276 objects without extending outward. It is hypothesized that the likelihood of triggering new  
277 convection is lower (since the surrounding environment is within smaller precipitation areas).  
278 Additionally, at this environmental scale, larger-scale environmental signals such as  
279 orographic cooling or valley winds may also be reflected in the cold pool signals. When the  
280 environmental scale is near the conventional value (89–100 km) (Figs. 5g and h), the cold  
281 pool intensity becomes weaker ( $<1 \text{ ms}^{-1}$ ) in the plain region and excessively strong ( $>5 \text{ ms}^{-1}$ )  
282 in the mountain region, which does not align with the precipitation patterns shows in Fig.

283 1.

284 To sum up, from the cold pool distribution under the environmental scale from 2–4 to  
285 89–100 km, we selected environmental scale in the range of 7–11 km as the appropriate  
286 environmental scale in southwest Taiwan.

287 To demonstrate that the cold pool distribution is reasonable based on the precipitation  
288 within our selected scales, we calculate the cold pool frequency distribution across  
289 ensemble members. From Fig. 6a, which is the cold pool calculated from the environmental  
290 scale in 7–11 km, the result shows that cold pool mainly occurred at the near-mountain  
291 region, where is the precipitation region. And in Fig. 6b, the cold pool frequency distribution  
292 using 89–100 km environmental scale shows that the cold pool mainly occurs in the coastal  
293 and high mountain regions, which is without obvious precipitation occurred. These high  
294 frequency regions of cold pool are due to the larger environmental and terrain cooling  
295 features and cause the bias. From the statistic of frequency, the cold pool calculated from  
296 the environmental scale in 7–11 km relatively better reflects the precipitation region. In  
297 addition, within the environmental scale of 7–11 km, the cold pool distribution not only  
298 concentrates near the mountainous region (500 m-1500 m altitude), which is close to the  
299 high-precipitation area shown in Fig. 4a, but also exhibits a clear line structure extending  
300 from the coastline to the inland. We infer that the cold pool signal results from different  
301 contributions to the negative buoyancy near the surface. We will further discuss the  
302 buoyancy contribution in the next subsection.



### 303 **3.3 Contribution for the negative buoyancy in cold pool**

304 The cold pool intensity in this study is calculated based on buoyancy, which is the  
 305 density potential temperature ( $\theta_v$ ) anomaly (Eqn.4). There are three terms contributing to  $\theta_v$   
 306 (Eqn.2): potential temperature, water vapor, and liquid water loading (i.e., cloud, ice, and  
 307 rainwater). To analyze the contribution from these three terms individually, we follow the  
 308 linearized buoyancy equation in Wilhelmson and Ogura (1972), and Grabowski and  
 309 Morrison (2021) below:

$$310 \quad B = g \left( \frac{\theta'}{\theta_0} + 0.608q'_v - q_l \right) \quad (\text{Eqn. 5}), \text{ where}$$

311  $\theta' = \theta - \tilde{\theta}$ ,  $q'_v = q_v - \tilde{q}_v$ , the tilde term ( $\sim$ ) is the spatial averaged with the environmental  
 312 scale domain. The  $\theta_0$  represents the reference state (Stevens, 2005, Yano et al., 2004).

313 From the linearized buoyancy equation, it shows the negative buoyancy contribution in cold  
 314 pools comes from negative potential temperature anomaly ( $\frac{\theta'}{\theta_0}$ ), reduced water vapor (0.608  
 315  $q'_v$ ), and increased liquid water loading ( $-q_l$ ). The negative potential temperature anomaly is  
 316 mainly from the evaporative cooling (roughly 2 order larger than radiative cooling in our study)  
 317 near the precipitation region under the environmental scale in 7–11 km. We use evaporative  
 318 cooling to refer to the effect caused by the negative potential temperature anomaly.

319 In Fig. 7, we calculate the cold pool contributions from these three terms for environmental  
 320 scale of 7.25 km. To discuss the spatial distribution of precipitation regions and the  
 321 contributions from the three terms to the cold pool and precipitation, we use three  
 322 precipitation thresholds to represent the regions with precipitation intensity above light

323 rainfall ( $5 \text{ mm h}^{-1}$ , red solid contour), moderate rainfall ( $15 \text{ mm h}^{-1}$ , red dashed contour), and  
324 heavy rainfall ( $50 \text{ mm h}^{-1}$ , green contour). The contribution to cold pool intensity from  
325 evaporative cooling (Figs.7a and b) are roughly similar to those from increased liquid water  
326 loading (Figs.7c and d), while the contribution from reduced water vapor (Figs.7e and f) is  
327 the smallest among the three; however, they have different spatial distribution over plain  
328 and mountain regions. Fig.7a shows that in the plain region, the cold pool contribution from  
329 evaporative cooling is mainly distributed around the regions above moderate rainfall  
330 intensity, while the cold pool intensity over the regions above light rainfall is relatively less  
331 pronounced. Fig. 7b shows that, in the mountainous region ( $X = 3\text{--}5 \text{ km}$ ), the cold pool  
332 caused by evaporative cooling is also primarily distributed around the moderate precipitation  
333 region. However, the cold pool intensity resulting from evaporative cooling in the  
334 mountainous region is weaker than that in the plain regions. A possible reason for this is the  
335 higher temperature in plains compared to mountains, which results in a higher saturation  
336 vapor pressure over plains. Consequently, the amount of water vapor available for  
337 evaporation is greater in plains than in mountainous areas. As a result, evaporative cooling  
338 is more effective in plains, leading to the formation of stronger cold pools compared to those  
339 in mountainous regions. Fig.7c shows that the cold pool caused by increased liquid water  
340 loading is primarily distributed within the center of the region above moderate rainfall  
341 intensity. In Fig.7d, in the vicinity of the mountain hill region, the cold pool caused by  
342 increased liquid water loading distribute within both the moderate precipitation region ( $X =$

343 3–5 km) and the heavy precipitation regions within the mountain hill area ( $X = 18\text{--}20$  km,  $Y$   
344  $= 8\text{--}15$  km, and  $X = 15$  km,  $Y = 22$  km). In Figs.7e and 7f, the distribution of the cold pool  
345 contribution from water vapor does not exhibit clear relationship with the precipitation  
346 regions either in the plain or mountainous areas, and its overall intensity remains relatively  
347 weak ( $<1$  m s<sup>-1</sup>). These results indicate that the cold pool associated with precipitation over  
348 Taiwan's plains and mountainous regions is primarily driven by evaporative cooling and  
349 increased liquid water loading.

350 To demonstrate the statistics of the cold pool from the three terms (evaporative cooling,  
351 liquid water loading, and water vapor) in Southwestern Taiwan across all ensemble  
352 members, we define the cold pool intensity greater than or equal to 1 m s<sup>-1</sup> as cold pool  
353 occurred. In Fig.8a, the result shows that the cold pool caused by evaporative cooling is  
354 mainly distributed in the mountain hill region and the plains near the mountain hill region.  
355 This may be due to the cold pool generated by evaporative cooling from convective  
356 precipitation, triggered when the sea breeze advances toward the mountain hill region in our  
357 simulation. The broader distribution of the cold pool in this area might be related to the  
358 widespread occurrence of convection in the plains near the mountainous region.

359 Next, the cold pool from the increased liquid water loading shows that the cold pool  
360 caused by increased liquid water loading is primarily distributed in the mountain hill and high  
361 mountain regions (Fig.8b). This distribution is associated with intense precipitation occurring  
362 in these mountainous areas. However, due to the smaller spatial extent of intense

363 precipitation, the frequency and spatial coverage of this cold pool are relatively smaller and  
364 less frequent compared to those caused by evaporative cooling.

365 Finally, the cold pool contributed from reduced water vapor is mainly distributed over  
366 the coastal plains (Fig.8c). And it roughly shows a distinct north-south boundary with the  
367 cold pool contributed from evaporative cooling (Fig.8a). The cold pool contributed from  
368 reduced water vapor may be due to local subsidence, in addition to the sea-breeze inflow,  
369 forming Horizontal Convective Rolls (HCR) (Dailey and Fovell, 1999) before the convective  
370 precipitation initiated from sea-breeze front and surface heating. This local subsidence  
371 brings relatively dry air with less water vapor, contributing to negative buoyancy in this region.  
372 While the cold pool contributions from water vapor induce some convection in the plain  
373 region, the intensity and occurrence are too small to be significant in our study.

374 Based on the results from different environmental scales and cold pool contributions,  
375 it is shown that the 7–11 km environmental scale can accurately reflect the cold pool  
376 distribution over Taiwan's complex topography. In Section 3.4, we evaluate the cold pool  
377 structure over this complex terrain by examining the evolution of the cold pool and  
378 convection in a cross-sectional view of the mountainous region.

### 379 **3.4 The evolution of cold pool and the convection interaction in mountain region**

380 To demonstrate the evolution and interaction between the cold pool and convection  
381 over Taiwan's complex topography, we evaluate the cold pool's development and its  
382 interaction with convection and near-surface inflow in the mountainous region. As indicated

383 by the cold pool evolution over a time interval from 14:40 to 15:00 LST (Figs. 9a to c), the  
384 cold pool and convection together induce lower boundary updrafts in less than 10 minutes,  
385 with the cold pool providing more convection triggering than the terrain. Additionally, the  
386 evolution shows that these updrafts (seen as new convection initiations) generate new cold  
387 pool signals, which can then induce updrafts in other locations.

388 We use vertical cross-section diagrams to illustrate the interactions among cold pools,  
389 convection, inflows, and terrain. These processes occur on a short convective timescale,  
390 with significant changes happening within approximately 20 minutes. In Fig. 9d (14:40), at  
391 around 10 km, the inflow collides with the leading edge of the cold pool near the mountain  
392 slope from the previous time step, triggering an updraft. This process exemplifies how the  
393 cold pool edge leads to new convection. Simultaneously, at around 18 km, a downburst  
394 generated by precipitation interacts with the outflow from the convection at 10 km. Ten  
395 minutes later (14:50, see Fig. 9e), convection and precipitation occur at around 15 km,  
396 accompanied by a negative  $\theta_v'$  signal at a height of 1–1.5 km, indicating cooling of about  
397 0.1–0.3 K. This process maintains the cold pool's presence below ( $X=12\text{--}15$  km).

398 From 14:50 to 15:00, precipitation shifts to steeper terrain at horizontal positions of 15–  
399 17 km. The accompanying cold pool accumulates along the mountain slope, inducing  
400 localized updrafts at 15 km and cold pool heights between 1.5 and more than 2 km, with  
401 vertical velocities reaching approximately  $1\text{ m s}^{-1}$ . These results demonstrate that the cold  
402 pools defined in this study over complex terrain exhibit preliminary evidence of interactions

403 among inflows, convection, cold pools, and terrain on small horizontal scales (approximately  
404 5 km) and short timescales (10–20 minutes). This provides a foundation for future research  
405 on the physical processes of convection and cold pools in complex terrain.

406

#### 407 **4. Implication of cold pool hotspot changes under global warming**

408 From the previous analysis, we identified an appropriate environmental scale (7–11 km)  
409 for the cold pool calculation over Taiwan's complex terrain. To highlight the advantages of  
410 our cold pool definition compared to the conventional scale (89–100 km), we used the cold  
411 pool defined by these two scales to analyze the afternoon thunderstorm simulations in  
412 TaiwanVVM. We then compared the changes in cold pool distribution between the current  
413 and PGW scenarios.

414 Based on Chen et al. (2024), where 30 ensemble members are simulated under local-  
415 circulation-dominated weather patterns during summer, the PGW scenario assumes an  
416 increase of 3 K in the entire atmospheric temperature (while maintaining the same relative  
417 humidity) compared to the current climate. In the current climate scenario, the precipitation  
418 hotspot is primarily concentrated in the mountainous region, as discussed in Section 3.1.  
419 Under the PGW scenario, the precipitation hotspot expands to the plains (Fig. 10a). The  
420 transition from the current climate to the PGW scenario shows the precipitation hotspot  
421 shifting from the mountainous region (Fig. 4a) to plain area near the mountains (Fig. 10b),  
422 with a frequency increase of up to 9%.

423 Under the PGW scenario, the cold pool calculated using the 7-11 km environmental scale  
424 shows a frequency distribution pattern (Fig. 10c) similar to that of the current climate (Fig.  
425 6a). The difference in cold pool frequency between the current and PGW scenarios (Fig.  
426 10d) shows an overall increase in frequency, with a greater concentration over the plain  
427 areas near the mountains compared to the mountainous regions. The cold pool frequency  
428 increases by approximately 9% in the corresponding regions. However, when the  
429 conventional environmental scale is used, the pattern of cold pool frequency distribution  
430 becomes less distinct (Fig. 10e), with differences of up to 20%, primarily reflecting cold pools  
431 driven by land, sea, and mountain-scale forcing (Fig. 10f). This preliminary diagnosis opens  
432 up future applications for studying convection–cold pool interactions over complex  
433 topography, which could contribute to longer and stronger convective lifetimes and  
434 intensities, as discussed in Chen et al., 2024.

## 435 **5. Summary and Conclusions**

436 This study investigates the identification of appropriate cold pool scales over Taiwan's  
437 complex topography, particularly during prevalent afternoon thunderstorms under a local-  
438 circulation dominated weather regime in summer. Using semi-realistic TaiwanVVM  
439 simulations covering the entire Taiwan, we identify the appropriate environmental scale for  
440 calculating cold pools. We hypothesize that the environmental scale should be large enough  
441 to cover the precipitation object while remaining small enough to avoid interference from  
442 land, sea, and mountain distributions. Our analysis identifies an optimal environmental scale

443 of approximately 7–11 km, corresponding to three times the precipitation object length  
444 between the 75th and 90th percentiles. This newly identified environmental scale better  
445 captures the cold pool frequency distribution in relation to precipitation hotspots, with cold  
446 pools characterized by evaporative cooling over plain area near the mountains and  
447 increased water loading over mountainous areas. We also demonstrate that cold pool  
448 interactions associated with topography can be clearly identified using this definition.

449 In conclusion, this study provides a new insight into identify the cold pool over complex  
450 topography in Taiwan through the environmental scale about 7–11 km based on the  
451 relationship with the precipitation object of afternoon thunderstorm. The selection of such an  
452 environmental scale for other tropical islands can follow our approach by first identifying the  
453 dominant precipitation object scale, then systematically increasing it to the conventional  
454 environmental scale to determine the optimal scale which further studies of cold pool related  
455 processes.

#### 456 **Data Availability Statement**

457 The data that support the findings of this study are available on request from the  
458 corresponding author.

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465 to conduct the simulation

466

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## References

468 Benjamin, T. B., 1968: Gravity currents and related phenomena. *Journal of Fluid*  
469 *Mechanics*, **31(2)**, 209–248

470 Bjorn Stevens, 2005: ATMOSPHERIC MOIST CONVECTION

471 <https://doi.org/10.1146/annurev.earth.33.092203.122658>

472 Byers, H. R., & Braham, R. R., 1949: The thunderstorm: Report of the thunderstorm  
473 project. Washington, DC: US Government Printing Office.

474 Chang, Y.-H., Chen, W.-T., Wu, C.-M., Moseley, C., and Wu, C.-C., 2021: Tracking the  
475 influence of cloud condensation nuclei on summer diurnal precipitating systems over  
476 complex topography in Taiwan, *Atmos. Chem. Phys.*, **21**, 16709–16725,  
477 <https://doi.org/10.5194/acp-21-16709-2021>

478 Chang Y-H, Chen W-T, Wu C-M, Kuo Y-H and Neelin, J-D, 2023: Identifying the deep553  
479 inflow mixing features in orographically-locked diurnal convection *Geophys Res.*  
480 *Lett.* 50 e2023GL103107

481 Chen, F. and Dudhia, J. 2001: Coupling an Advanced Land Surface–Hydrology Model with  
482 the Penn State–NCAR MM5 Modeling System. Part I: Model Implementation and

- 483 Sensitivity, *Mon. Weather Rev.*, **129**, 569–585, doi:10.1175/1520-  
484 0493(2001)129<0569:Caalsh>2.0.Co;2.
- 485 Chen, F., Mitchell, K., Schaake, J., Xue, Y., Pan, H.-L., Koren, V., Duan, Q. Y., Ek, M., and  
486 Betts, A., 1996: Modeling of land surface evaporation by four schemes and comparison  
487 with FIFE observations, *J. Geophys. Res.-Atmos.*, **101**, 7251–  
488 7268, doi:10.1029/95jd02165.
- 489 Chen, P.-J., Chen, W.-T., Wu, C.-M., & Yo, T.-S., 2021: Convective cloud regimes from a  
490 classification of object-based CloudSat observations over Asian-Australian monsoon  
491 areas. *Geophysical Research Letters*, **48(10)**, e2021GL092733.  
492 <https://doi.org/10.1029/2021GL092733>
- 493 Chen, W.-T., Y.-H. Chang, C.-M. Wu, H-Y. Huang, 2024: The future extreme precipitation  
494 systems of orographically locked diurnal convection: The benefits of using large-eddy  
495 simulation ensembles, *Environmental Research: Climate*. [https://doi.org/10.1088/2752-  
496 5295/ad557d](https://doi.org/10.1088/2752-5295/ad557d)
- 497 Chen, X., Zhao, K., Xue, M., Zhou, B., Huang, X., & Xu, W., 2015: Radar-observed diurnal  
498 cycle and propagation of convection over the Pearl River Delta during Mei-Yu season.  
499 *Journal of Geophysical Research: Atmospheres*, **120(24)**, 12557–12575.  
500 doi:10.1002/2015JD023872
- 501 Dailey, P.S., Fovell, R.G., 1999: Numerical simulation of the interaction between the sea-  
502 breeze front and horizontal convective rolls. Part I: offshore ambient flow. *Mon. Weather*

- 503 Rev. **127 (5)**, 858–878. <https://doi.org/10.1175/1520->
- 504 [0493\(1999\)127<0858:nsotib>2.0.co;2](https://doi.org/10.1175/1520-0493(1999)127<0858:nsotib>2.0.co;2).
- 505 Doswell, C. A., 2001: Severe convective storms—An overview. In C. A. Doswell (Ed.),
- 506 Severe convective storms, Meteorological Monographs. Boston, MA: American
- 507 Meteorological Society.
- 508 Drager, A. J., & van den Heever, S. C., 2017: Characterizing convective cold pools.
- 509 Journal of Advances in Modeling Earth Systems, **9**, 1091–1115.
- 510 <https://doi.org/10.1002/2016MS000788>
- 511 Droegemeier, K. K., & Wilhelmson, R. B., 1985: Three-dimensional numerical modeling of
- 512 convection produced by interacting thunderstorm outflows. Part I: Control simulation and
- 513 low-level moisture variations. Journal of the Atmospheric Sciences, **42(22)**, 2381–2403.
- 514 Feng, Z., Hagos, S., Rowe, A. K., Burleyson, C. D., Martini, M. N., & de Szoeke, S. P.,
- 515 2015: Mechanisms of convective cloud organization by cold pools over tropical warm
- 516 ocean during the AMIE/DYNAMO field campaign. Journal of Advances in Modeling
- 517 Earth Systems, **7**, 357–381. doi:10.1002/2014MS000384
- 518 Fournier, M. B., & Haerter, J. O., 2019: Tracking the gust fronts of convective cold pools.
- 519 Journal of Geophysical Research, **124**, 11,103–11,117.
- 520 <https://doi.org/10.1029/2019JD030980>
- 521 Henneberg, O., Meyer, B., & Haerter, J. O., 2020: Particle-based tracking of cold pool gust
- 522 fronts. Journal of Advances in Modeling Earth Systems, **12**, e2019MS001910.

523 <https://doi.org/10.1029/2019MS001910>

524 Jian, H.-W., Chen, W.-T., Chen, P.-J., Wu, C.-M., & Rasmussen, K. L., 2021: The

525 synoptically-influenced extreme precipitation systems over Asian-Australian monsoon

526 region observed by TRMM precipitation radar. *Journal of the Meteorological Society of*

527 *Japan*, **99(2)**, 269–285. <https://doi.org/10.2151/jmsj.2021-013>

528 Jun-Ichi Yano, Jean-Pierre Chaboureau, Françoise Guichard, 2005A: A generalization of

529 CAPE into potential-energy convertibility. <https://doi.org/10.1256/qj.03.188>

530 Keulegan, G., 1957: An experimental study of the motion of saline water from locks into

531 fresh water channels (Technical Report 5168): National Bureau of Standards.

532 Krishna, U. V. M., Das, S. K., Deshpande, S. M., & Pandithurai, G., 2021: Physical

533 processes controlling the diurnal cycle of convective storms in the Western Ghats.

534 *Scientific Reports*, **11(1)**, 14103. doi:[10.1038/s41598-021-93173-0](https://doi.org/10.1038/s41598-021-93173-0)

535 Kuenen, P. H., 1950: Turbidity currents of high density. 18th International Geological

536 Congress (1948) (Vol. 8). London. Reports, pt.

537 Kuo, K.-T., and C.-M. Wu, 2019: The precipitation hotspots of afternoon thunderstorms

538 over the Taipei basin: Idealized numerical simulations. *J. Meteor. Soc. Japan*, **97**, 501–

539 517.

540 Kuo, Y.-H., & Neelin, J. D., 2022: Conditions for convective deep inflow. *Geophysical*

541 *Research Letters*, **49**, e2022GL100552. <https://doi.org/10.1029/2022GL100552>

542 Miao, J.-E., and M.-J. Yang, 2020: A Modeling Study of the Severe Afternoon

- 543 Thunderstorm Event at Taipei on 14 June 2015: The Roles of Sea Breeze,  
544 Microphysics, and Terrain. *J. Meteor. Soc. Japan*, **98**, 129-152.
- 545 Miao, J.-E., and M.-J. Yang, 2022: The Impacts of Midlevel Moisture on the Structure,  
546 Evolution, and Precipitation of Afternoon Thunderstorms: A Real-Case Modeling Study  
547 at Taipei on 14 June 2015. *J. Atmos. Sci.*, **79**, 1837–1857.
- 548 Nesbitt S-W and Zipser E-J, 2003: The diurnal cycle of rainfall and convective intensity  
549 according to three years of TRMM measurements *J. Clim.* **16** 1456-1475
- 550 Romatschke, U., & Houze, R. A., 2011: Characteristics of precipitating convective systems  
551 in the South Asian monsoon. *Journal of Hydrometeorology*, **12(1)**, 3–26.  
552 <https://doi.org/10.1175/2010JHM1289.1>
- 553 Schlemmer, L. and Hohenegger, C., 2015: Modifications of the atmospheric moisture field  
554 as a result of cold-pool dynamics, *Q. J. Roy. Meteor. Soc.*, **142**, 30–42,  
555 <https://doi.org/10.1002/qj.2625>.
- 556 Song, Z., & Zhang, J., 2020: Diurnal variations of summer precipitation linking to the  
557 topographical conditions over the Beijing-Tianjin-Hebei region. *Scientific Reports*, **10(1)**,  
558 9701. <https://doi.org/10.1038/s41598-020-65743-1>
- 559 Thorpe, A., Miller, M., & Moncrieff, M., 1982: Two-dimensional convection in non-constant  
560 shear: A model of mid-latitude squall lines. *Quarterly Journal of The Royal  
561 Meteorological Society*, **108(458)**, 739–762.
- 562 Tompkins, A. M., 2001a: Organization of tropical convection in low vertical wind shears:

563 The role of cold pools. *Journal of the Atmospheric Sciences*, **58(13)**, 1650–1672.

564 Tompkins, A. M. (2001b). Organization of tropical convection in low vertical wind shears:

565 The role of water vapor. *Journal of the atmospheric sciences*, **58(6)**, 529–545.

566 Tsai, W.-M., and C.-M. Wu, 2017: The environment of aggregated deep convection, *J.*

567 *Adv. Model. Earth Syst.*, **9**, 2061–2078, doi:10.1002/2017MS000967.

568 Von Karman, T., 1940: The engineer grapples with nonlinear problems. *Bulletin of the*

569 *American Mathematical Society*, **46(8)**, 615–683.

570 Wilhelmson, R., and Y. Ogura, 1972: The Pressure Perturbation and the Numerical

571 Modeling of a Cloud. *J. Atmos. Sci.*, **29**, 1295–1307, <https://doi.org/10.1175/1520->

572 [0469\(1972\)029<1295:TPPATN>2.0.CO;2](https://doi.org/10.1175/1520-0469(1972)029<1295:TPPATN>2.0.CO;2).

573 Webster, P. J., & Lukas, R., 1992: TOGA COARE: The coupled ocean–atmosphere

574 response experiment. *Bulletin of the American Meteorological Society*, **73(9)**, 1377–

575 1416.

576 Wojciech W. Grabowski and Hugh Morrison, 2021: Supersaturation, buoyancy, and deep

577 convection dynamics <https://acp.copernicus.org/articles/21/13997/2021/#section13>

578 Wu C-M and Arakawa A, 2011: Inclusion of surface topography into the vector vorticity

579 equation model (VVM) *J. Adv. Model. Earth Syst.* **3** M04002

580 Wu C-M, Lin H-C, Cheng F-Y and Chien M-H, 2019: Implementation of the land surface

581 processes into a vector vorticity equation model (VVM) to study its impact on afternoon

582 thunderstorms over complex topography in Taiwan Asia-Pac. *J. Atmos. Sci.* **55** 701-717

583 Yang G-Y and Slingo J, 2001: The diurnal cycle in the tropics *Mon. Weather Rev.* **129** 725  
584 784-801

585 Zuidema, P., Torri, G., Muller, C., & Chandra, A., 2017: A survey of precipitation-induced  
586 atmospheric cold pools over oceans and their interactions with the larger-scale  
587 environment. *Surveys in Geophysics*, **38(6)**, 1283–1305.

588 List of Figures

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593 Fig. 2 The initial profiles of equivalent potential temperature and saturated equivalent  
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596 potential temperature, and green lines indicate saturated equivalent potential temperature.  
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599 are both 960hPa, level of free convection (LFC) are 820 and 840hPa. The convective  
600 inhibition (CIN) are 55 and 45  $\text{J kg}^{-1}$ , and the convective available potential energy (CAPE)  
601 are 1600 and 3200  $\text{J kg}^{-1}$ , respectively.

602 Fig. 3 Schematic diagram of the precipitation object length. (a) shows a color-shaded

603 snapshot of the precipitation rate in the mountainous region for one ensemble member  
604 during the afternoon, while (b) shows the corresponding snapshot of the precipitation  
605 object length at the same time. The grey shading represents the terrain, with black  
606 contours indicating altitudes of 0, 500, and 1500 meters. The edge length of the black box  
607 represents the environmental scale (approximately 10 km), based on three times the 75<sup>th</sup>  
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 644 LST. Frequency is defined by the number of occurrences where cold pool intensity is  $\geq 1$   
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 649 pool intensity ( $\text{m s}^{-1}$ ), and the red contour marks the updraft region where vertical velocity  
 650 is  $1 \text{ m s}^{-1}$  at a 1 km height. In the lower panel (d–f), the vertical cross-section follows the  
 651 black dashed line in (a–c), oriented from Northwest to Southeast. In the upper portion of  
 652 (d–f), green bars represent terrain height (km), blue line indicates cold pool height (km),  
 653 purple shading denotes the negative  $\theta'_v$ , which is  $(\theta_v - \tilde{\theta}_v)$  (K). Black contour  
 654 represents where the combined cloud water and ice water mixing ratio higher than  $0.01 \text{ g}$   
 655  $\text{kg}^{-1}$ , and bright shading highlights rainwater regions where the mixing ratio exceeds  $0.5$   
 656  $\text{g kg}^{-1}$ . Purple vectors represent the horizontal wind anomaly, calculated by subtracting  
 657 the horizontal wind at  $X = 0 \text{ km}$  (sea-breeze inflow) from the original wind field, and are  
 658 projected along the dashed line. These vectors are combined with vertical wind ( $\text{m s}^{-1}$ ),  
 659 while red vectors highlight vertical winds exceeding  $1 \text{ m s}^{-1}$ . The red contour represents  
 660 a region where the updraft velocity is higher than  $3 \text{ m s}^{-1}$ . In the lower portion, blue bars  
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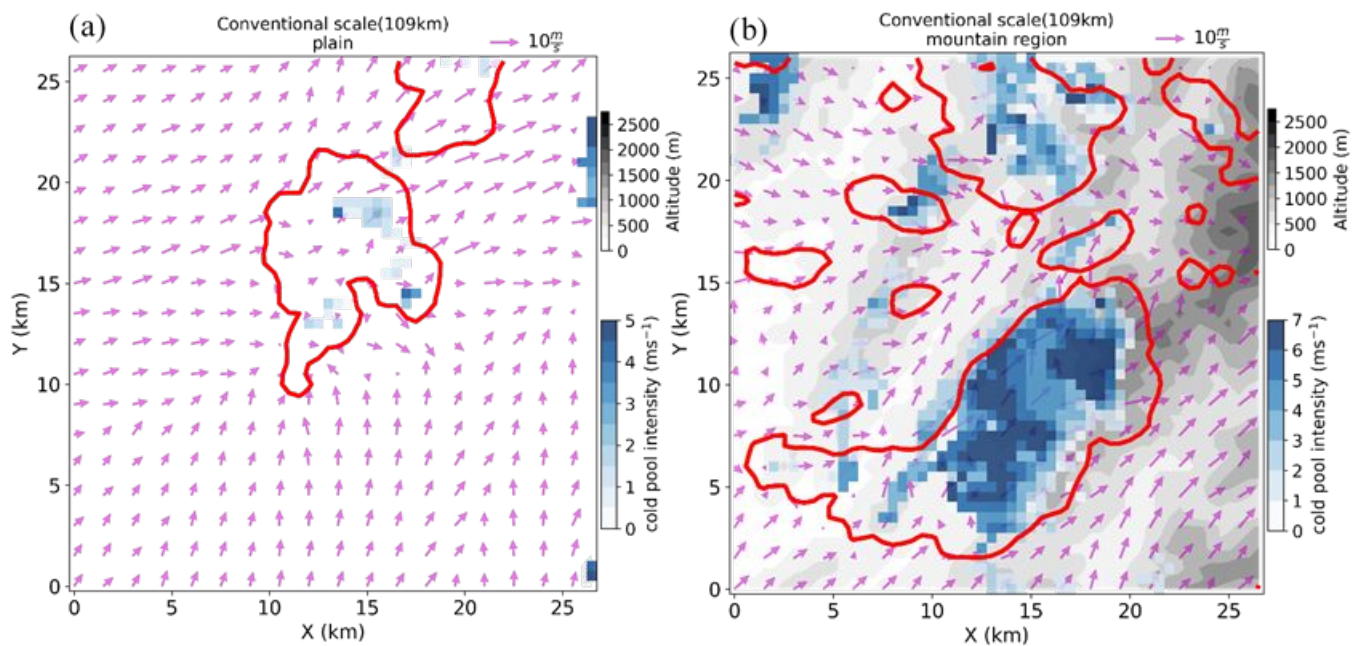
663 cold pools under environmental scales of 7–11 km (c, d) and 89–100 km (e, f) in the PGW  
 664 scenario, along with the difference between the two scenarios (PGW - current). Frequency  
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 666 contours represent altitudes at 0, 500, and 1500 meters, respectively.

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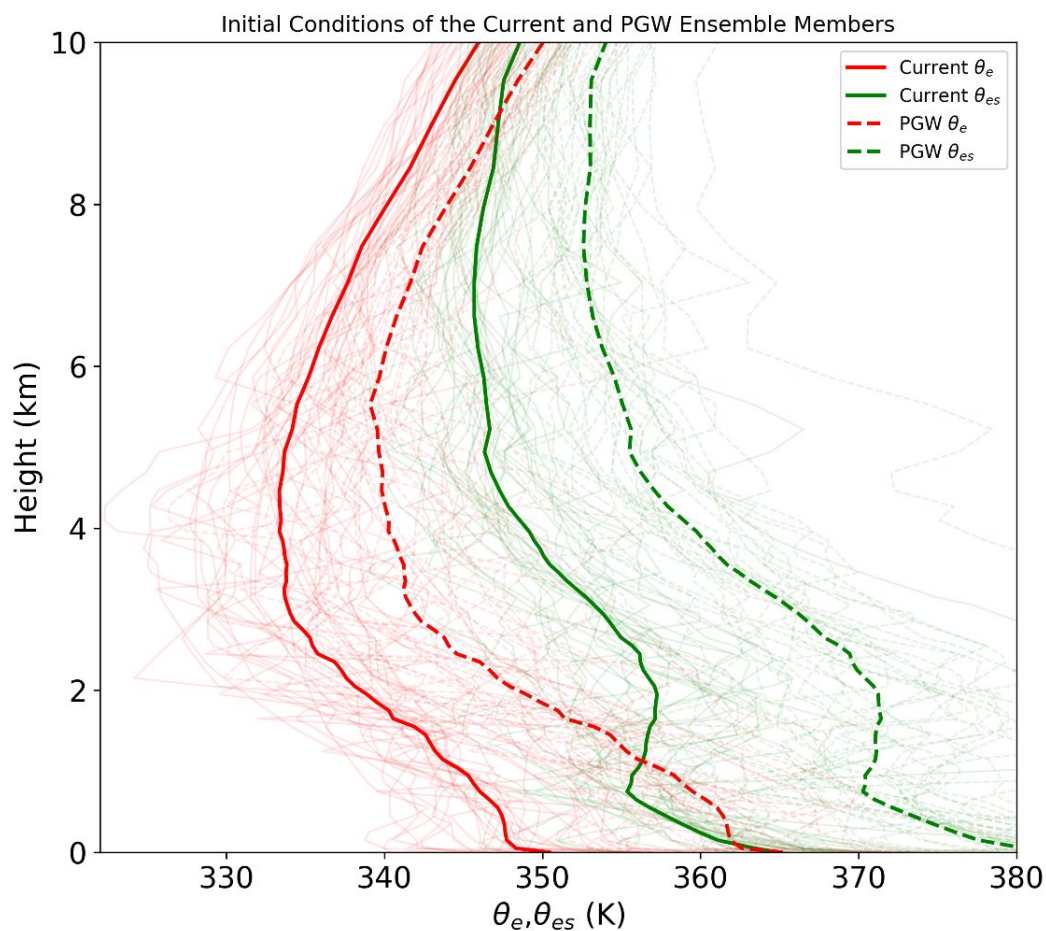
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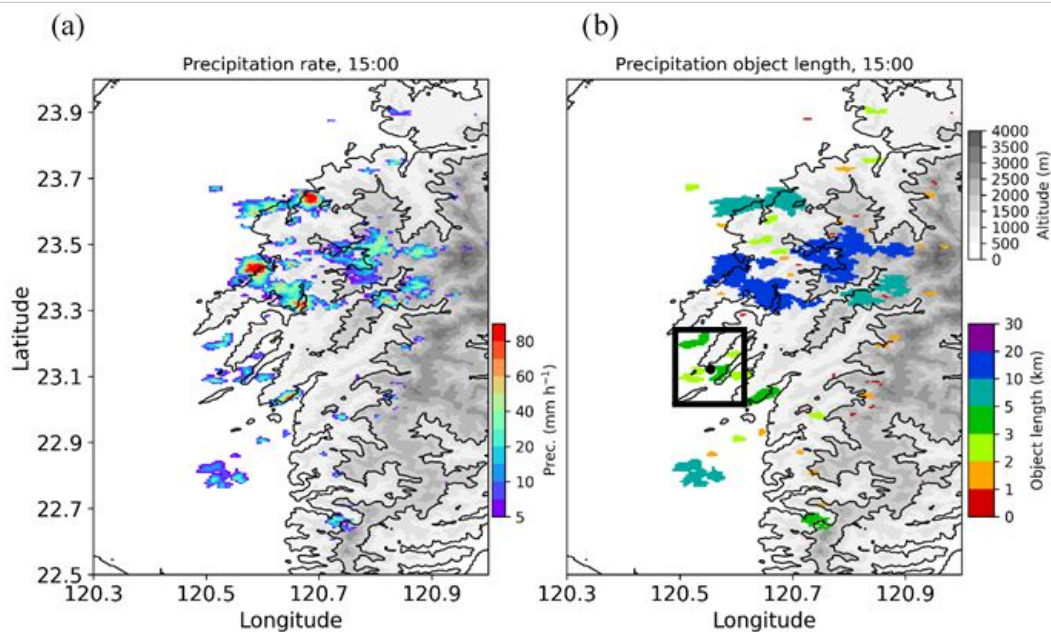
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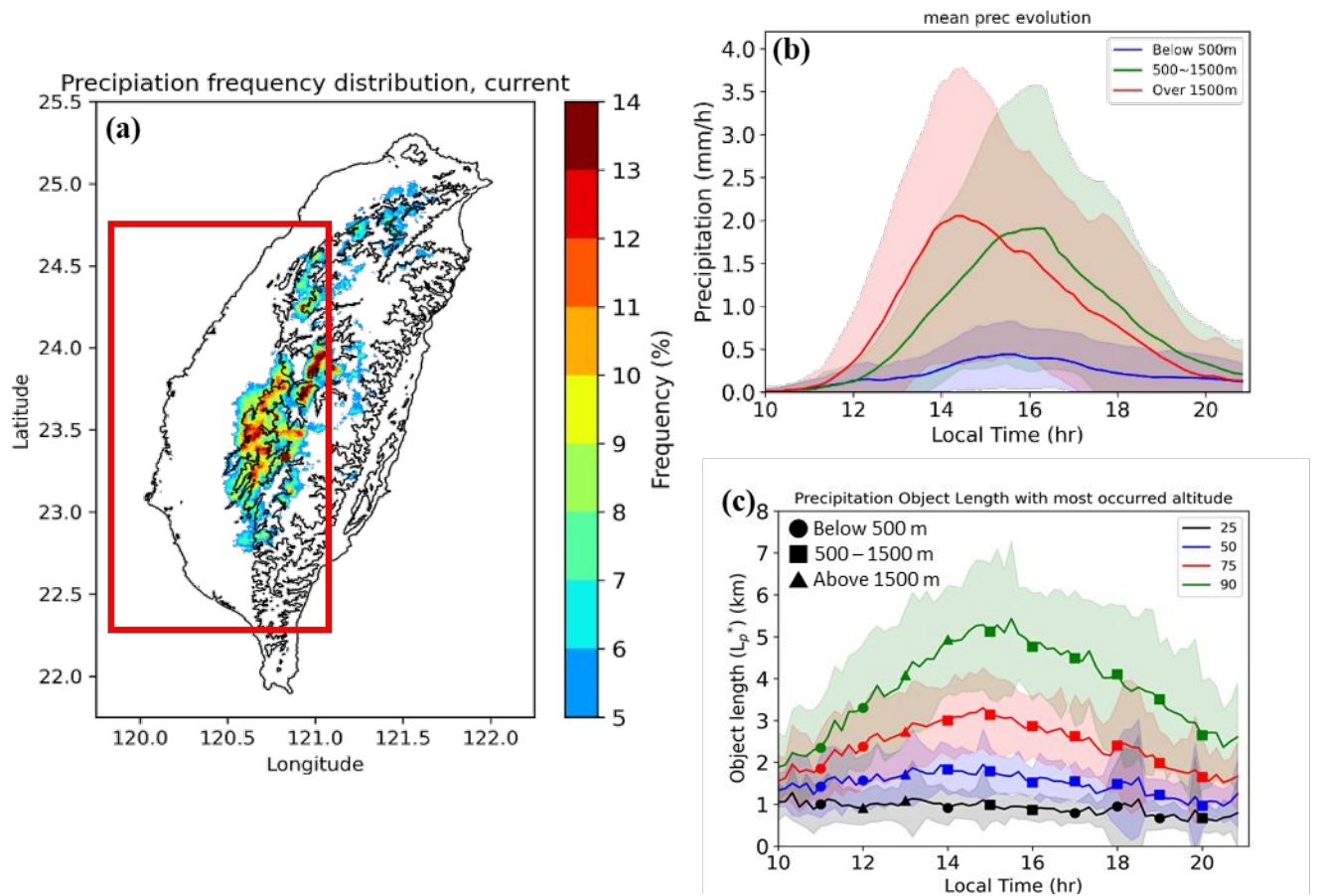
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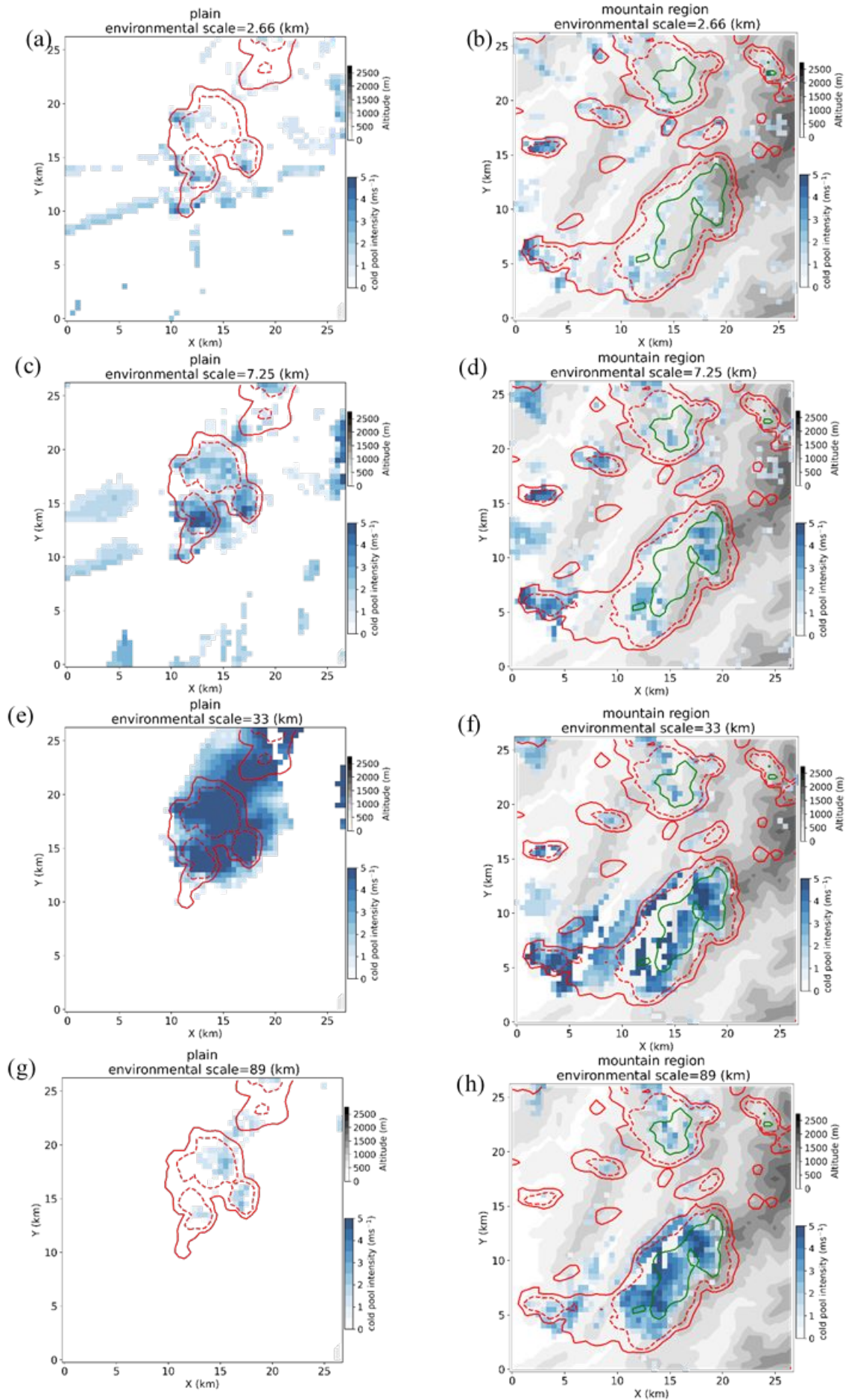


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701 Fig. 4 (a) The frequency (%) of precipitation exceeding  $5 \text{ mm h}^{-1}$  (shaded), based on  
 702 statistics from 14 to 21 LST across 30 ensemble members. Black contours represent  
 703 altitudes of 0, 500, and 1500 meters. The red box highlights the study region with relatively  
 704 high precipitation frequency. (b) The precipitation evolution in three sea-level height  
 705 regions (below 500 m, 500–1500 m, and above 1500 m) within the red box from (a). The  
 706 line represents the mean, and the shaded area indicates one standard deviation of  
 707 precipitation across 30 ensemble members. (c) The evolution of precipitation object length  
 708 over four percentiles (25<sup>th</sup>, 50<sup>th</sup>, 75<sup>th</sup>, 90<sup>th</sup>) per time period for each ensemble member.  
 709 The line represents the mean, and the shaded area indicates one standard deviation of  
 710 precipitation object length across 30 ensemble members. The markers indicate different  
 711 sea-level height regions where the mode of precipitation object centers falls within the four  
 712 percentiles across all 30 ensemble members. Dots represent areas below 500 m (plain),  
 713 squares correspond to regions between 500 and 1500 m (mountain hill region), and

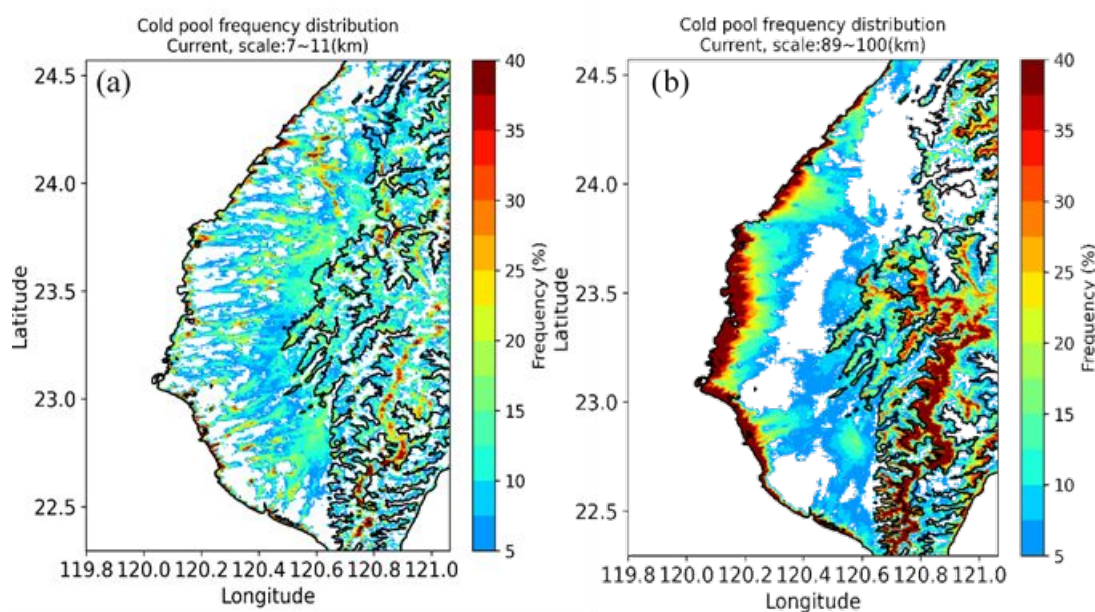
714 triangles denote areas above 1500 m (high-mountain region).

715





717 Fig. 5 The precipitation at rates of 5, 15, and 50 mm h<sup>-1</sup> (red solid, red dashed, and green  
 718 solid contours, respectively) and cold pool intensity (blue shading, unit: m s<sup>-1</sup>) calculated  
 719 using different environmental scales in the plain (a, c, e, g) and mountainous (b, d, f, h)  
 720 regions for one ensemble member. Grey shading represents sea-level height (in meters).  
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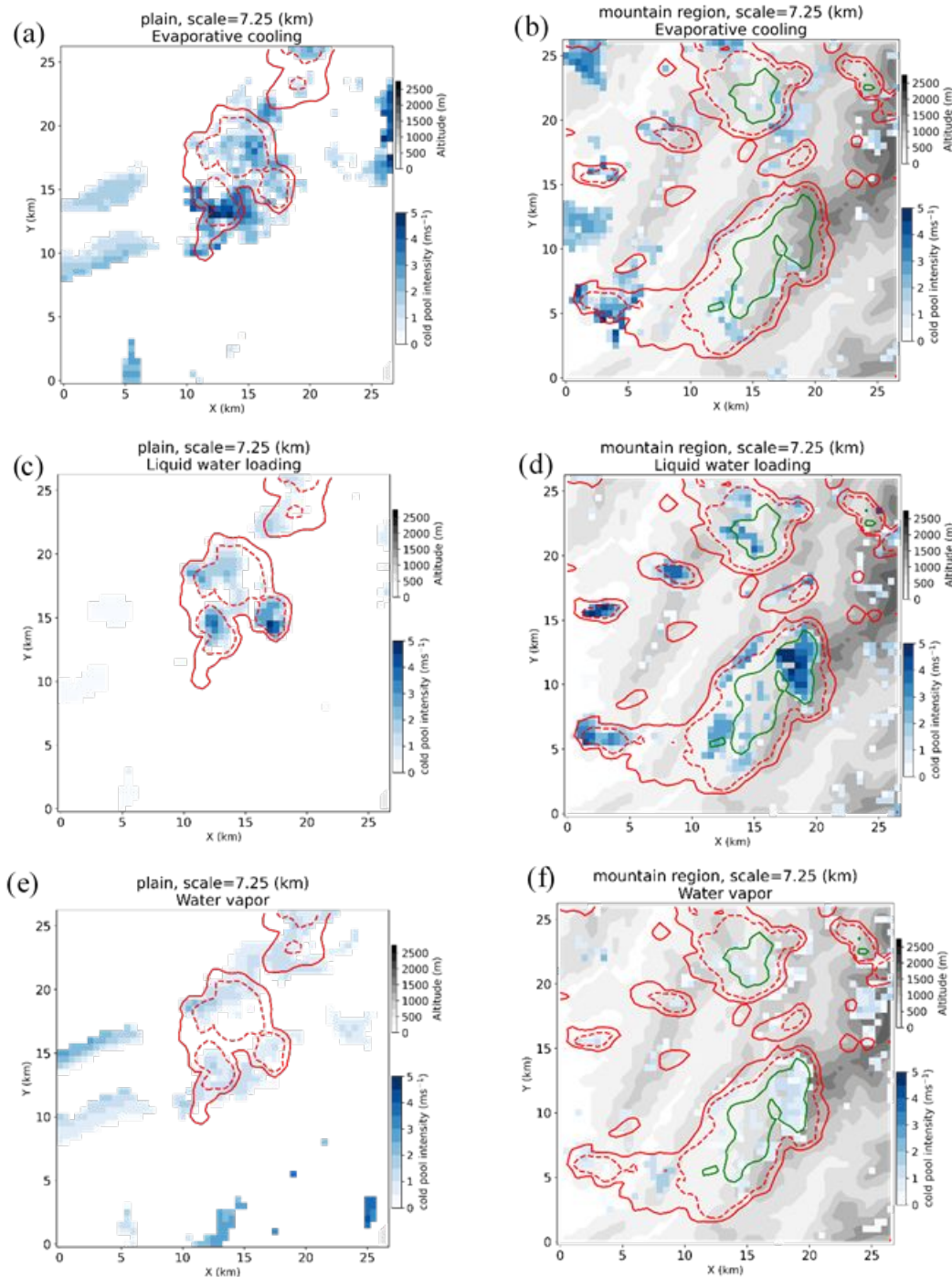


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723 Fig. 6 The occurrence frequency map of cold pools under the appropriate environmental  
 724 scale of (a) 7–11 km and (b) 89–100 km during 14–21 LST. Frequency is defined by the  
 725 number of occurrences where cold pool intensity is  $\geq 1$  m s<sup>-1</sup>. Contours represent altitudes  
 726 of 0, 500, and 1500 meters.

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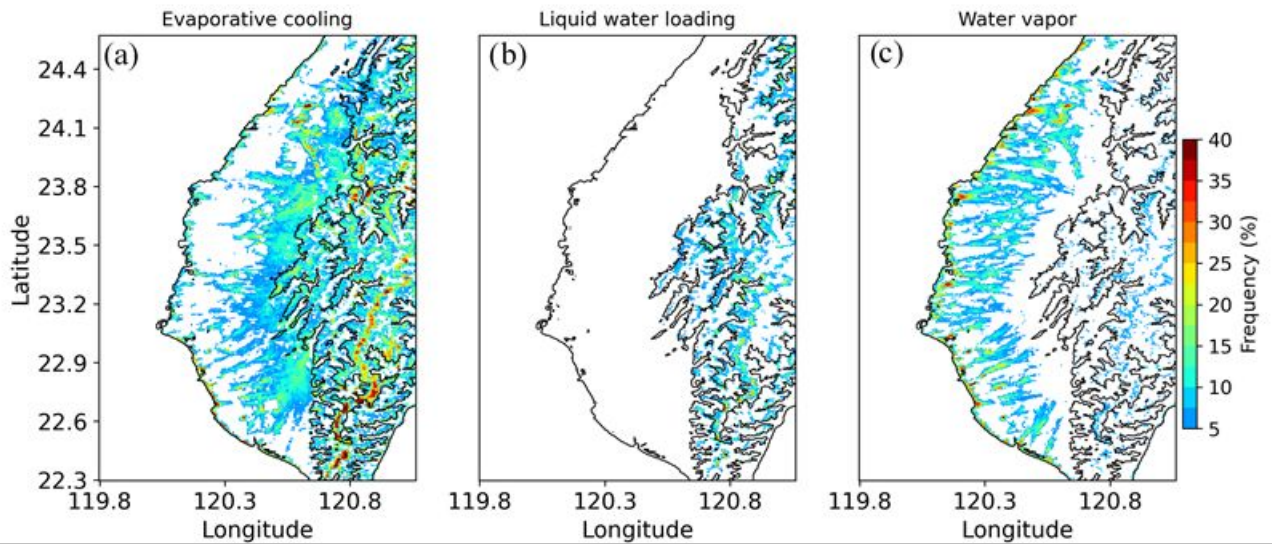
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730 Fig. 7 The negative buoyancy contribution to the cold pool, under the appropriate  
 731 environmental scale (7–11 km, represented here by a scale of 7.25 km), is shown in Figs.  
 732 5c and 5d for the plain (a, c, e) and mountainous (b, d, f) regions. Panels (a, b), (c, d), and  
 733 (e, f) correspond to the contributions from evaporative cooling, increased liquid water  
 734 content loading, and reduced water vapor effects, respectively. Precipitation rates of 5,  
 735 15, and 50 mm h<sup>-1</sup> are shown with red solid, red dashed, and green solid contours,

736 respectively, while cold pool intensity is shown in blue shading (unit:  $\text{m s}^{-1}$ ). Grey shading  
 737 represents altitude (in meters).

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741 Fig. 8 The occurrence frequency map of cold pools driven by negative buoyancy  
 742 contributions from (a) evaporative cooling, (b) increased liquid water loading, and (c)  
 743 reduced water vapor under the appropriate environmental scale (7–11 km) during 14–21  
 744 LST. Frequency is defined by the number of occurrences where cold pool intensity is  $\geq 1$   
 745  $\text{m s}^{-1}$ . Contours represent altitudes of 0, 500, and 1500 meters, respectively.

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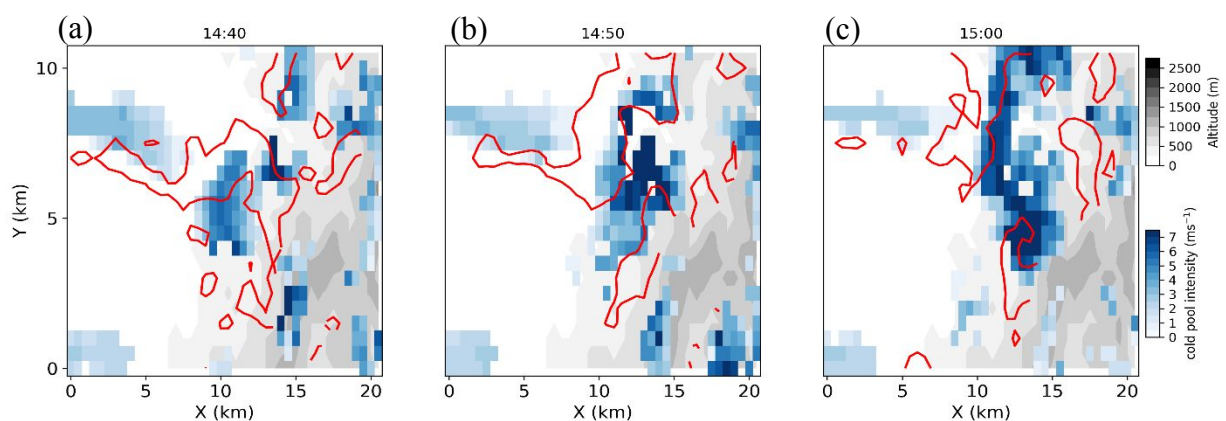
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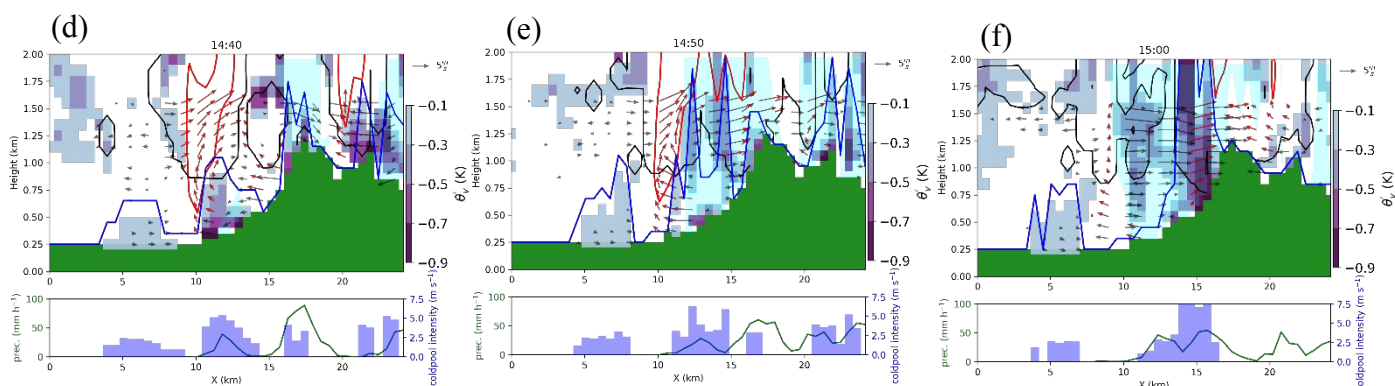
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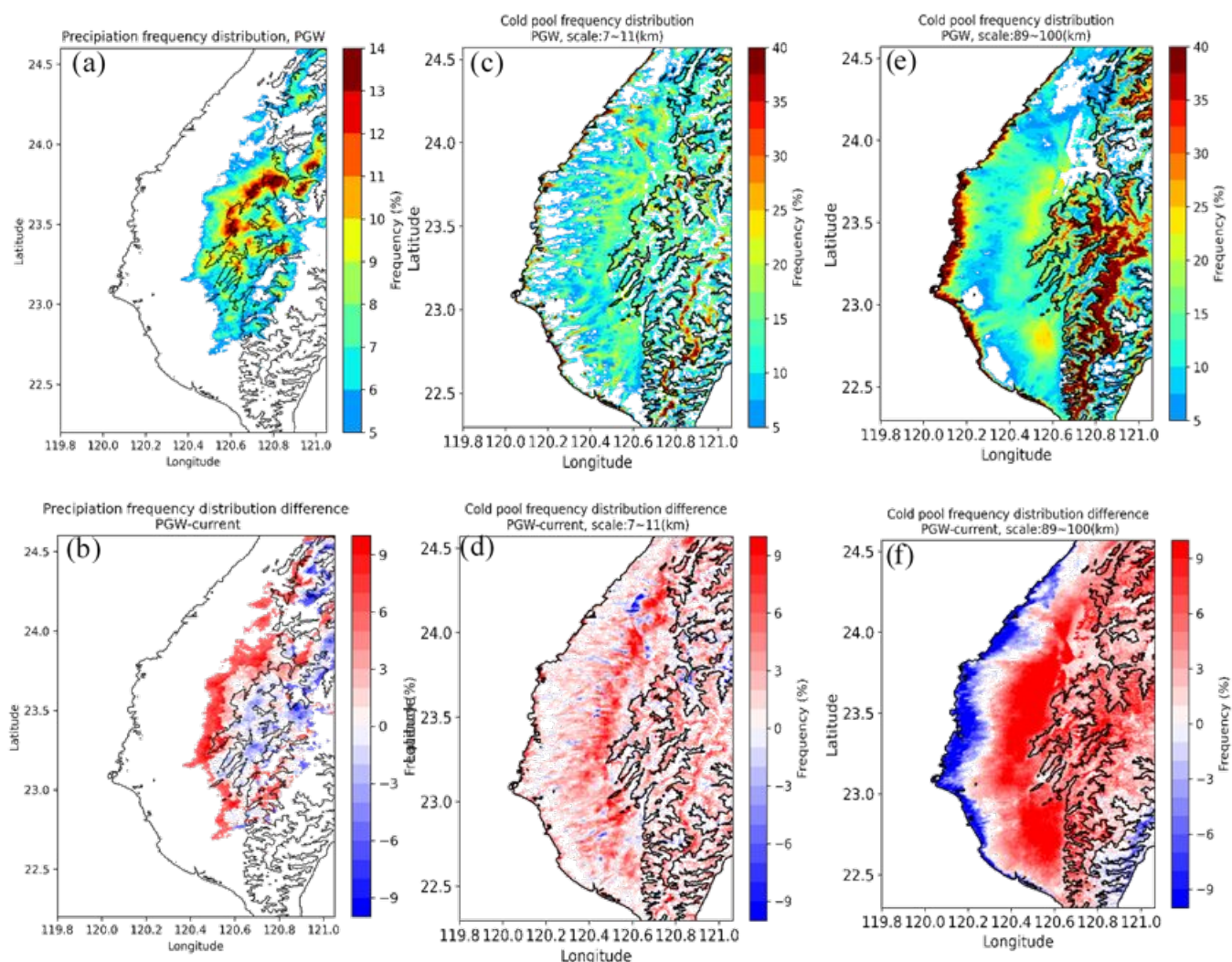
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Fig. 9 The evolution of the horizontal (a–c) and vertical cross-section (d–f) from 14:40 to 15:00 LST with the cold pool calculated using the 7–11 km environmental scale. In the upper panel (a–c), grey shading represents sea-level height (m), blue shading shows cold pool intensity ( $\text{m s}^{-1}$ ), and the red contour marks the updraft region where vertical velocity is  $1 \text{ m s}^{-1}$  at a 1 km height. In the lower panel (d–f), the vertical cross-section follows the black dashed line in (a–c), oriented from Northwest to Southeast. In the upper portion of (d–f), green bars represent terrain height (km), blue line indicates cold pool height (km), purple shading denotes the negative  $\theta'_v$ , which is  $(\theta_v - \tilde{\theta}_v)$  (K). Black contour represents where the combined cloud water and ice water mixing ratio higher than  $0.01 \text{ g kg}^{-1}$ , and bright shading highlights rainwater regions where the mixing ratio exceeds  $0.5 \text{ g kg}^{-1}$ . Purple vectors represent the horizontal wind anomaly, calculated by subtracting the horizontal wind at  $X = 0 \text{ km}$  (sea-breeze inflow) from the original wind field, and are projected along the dashed line. These vectors are combined with vertical wind ( $\text{m s}^{-1}$ ), while red vectors highlight vertical winds exceeding  $1 \text{ m s}^{-1}$ . The red contour represents a region where the updraft velocity is higher than  $3 \text{ m s}^{-1}$ . In the lower portion, blue bars depict cold pool intensity ( $\text{m s}^{-1}$ ), and the green line represents precipitation rate ( $\text{mm h}^{-1}$ ).



780

781 Fig. 10 The occurrence frequency map from 14:00 to 21:00 LST for precipitation (a, b) and

782 cold pools under environmental scales of 7–11 km (c, d) and 89–100 km (e, f) in the

783 PGW scenario, along with the difference between the two scenarios (PGW - current).

784 Frequency is defined by the number of occurrences where cold pool intensity is  $\geq 1$  m785  $s^{-1}$ . Black contours represent altitudes at 0, 500, and 1500 meters, respectively.

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## 792 List of Tables

793

794 Table 1 Configuration of the TaiwanVVM for semi-realistic simulations under current and  
 795 Pseudo Global Warming (PGW) scenarios.

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797

798 Table 1 Configuration of the TaiwanVVM for semi-realistic simulations under current and  
 799 Pseudo Global Warming (PGW) scenarios.

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Domain size (grid points)	512 km * 512 km * 19260 m 1024 * 1024 * 70 grids
Horizontal resolution	500 m
Vertical resolution	100 m (< 3900 m) Stretch up to 955 m (3900 m ~ model top)
Time step	10 secs
Simulation duration	24 h (00:00 ~ 24:00)
Lateral boundary condition	Double periodic
Output frequency	10 mins
PGW scenario	Temperature + 3K Keeps the same Relative Humidity

## **Identifying Cold Pool scales over Complex Topography using TaiwanVVM simulations**

We sincerely appreciate the editors' careful evaluation and constructive suggestions. In response to the editors' feedback, all changes have been accepted, and the final version is clean without any revision marks. Additionally, we have made a comprehensive revision to clarify our approach to defining the environmental scale over Taiwan's complex topography. Additionally, we have refined the manuscript to ensure clarity, incorporating input from English editing experts.